

Existence of Electron spin

(1)

The Dirac's Hamiltonian in a stationary central field represented by potential $V(r)$ is

$$H = c \vec{\alpha} \cdot \vec{P} + \beta mc^2 + V(r) \quad \text{--- (1)}$$

For a particle under the influence of a central force the torque is zero.

Therefore the angular momentum is a constant of

motion in central field.

The operator commutes with H is a constant of motion. Let us examine this for x -component of momentum.

$$\begin{aligned} i\hbar \frac{dL_x}{dt} &= [L_x, H] = [L_x, c \vec{\alpha} \cdot \vec{P} + \beta mc^2 + V(r)] \\ &= [L_x, c \{ \alpha_x P_x + \alpha_y P_y + \alpha_z P_z \} + \beta mc^2 + V(r)] \\ &= [L_x, c \alpha_x P_x] + [L_x, c \alpha_y P_y] + [L_x, c \alpha_z P_z] \\ &\quad + [L_x, \beta mc^2] + [L_x, V(r)] \quad \text{--- (2)} \end{aligned}$$

But L_x commute with the term P_y, P_z

$$[L_x, c\alpha_x p_x] = [L_x, \beta m c^2] = [L_x, V(r)] = 0 \quad (2)$$

Therefore, we get

$$\begin{aligned} i\hbar \frac{dL_x}{dt} &= [L_x, c\alpha_y p_y] + [L_x, c\alpha_x p_z] - (3) \\ &= [L_x, c\alpha_y p_y] = c\alpha_y [L_x, p_y] \\ &= c\alpha_y [y p_z - z p_y, p_y] \\ &= c\alpha_y \{ [y p_z, p_y] - [z p_y, p_y] \} \\ &= c\alpha_y [[y, p_y] p_z - y [p_z, p_y] - 0] \\ &= c\alpha_y [i\hbar p_z - 0] = c\alpha_y i\hbar p_z - (4) \end{aligned}$$

We can also calculate

$$[L_x, c\alpha_z p_z] = c\alpha_z (-i\hbar p_y) - (5)$$

$$\begin{aligned} i\hbar \frac{dL_x}{dt} &= c\alpha_y (i\hbar p_z) + c\alpha_z (-i\hbar p_y) \\ &= -i\hbar c (\alpha_z p_y - \alpha_y p_z) \\ &\neq 0 \end{aligned} \quad - (6)$$

$$\frac{dL_x}{dt} \neq 0, L_x \neq \text{constant}$$

Hence according to Dirac's theory the x -component⁽³⁾ of orbital angular momentum moving in a central electrostatic field is not a constant of motion. So we can find another operator such that the commutator of its x -component with H is equal and opposite of the right side of 5. So that the sum of this operator and L is then a constant of motion which can be interpreted as the total angular momentum.

It is not difficult to see that the desire operator is multiple of $\vec{\sigma}^1$,

$$\vec{\sigma}^1 = \begin{bmatrix} \vec{\sigma} & 0 \\ 0 & \vec{\sigma} \end{bmatrix}$$

Let us take x -component of $\vec{\sigma}^1$,

$$i\hbar \frac{d\sigma x^1}{dt} = [\sigma x^1, H]$$

$$\vec{\sigma}_{x^1} = \begin{bmatrix} \sigma x & 0 \\ 0 & \sigma x \end{bmatrix}$$

$$\text{i.e. } i\hbar \frac{d\sigma x^1}{dt} = \left[\sigma x^1, (c\vec{\alpha} \cdot \vec{p} + \beta mc^2 + V(r)) \right]$$

$$= \left[\sigma x^1, c(\alpha_x p_x + \alpha_y p_y + \alpha_z p_z) + \beta mc^2 + V(r) \right] - (7)$$

But σx^1 commutes with every quantity in above except α_y and α_z

$$\sigma x' \beta - \beta \sigma x' = \begin{bmatrix} \sigma x & 0 \\ 0 & \sigma x \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} - \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} \sigma x & 0 \\ 0 & \sigma x \end{bmatrix} \quad (4)$$

$$= \begin{bmatrix} \sigma x & 0 \\ 0 & -\sigma x \end{bmatrix} - \begin{bmatrix} \sigma x & 0 \\ 0 & -\sigma x \end{bmatrix} = 0$$

$$\sigma x' \alpha_x - \alpha_x \sigma x' = \begin{bmatrix} \sigma x & 0 \\ 0 & \sigma x \end{bmatrix} \begin{bmatrix} 0 & \sigma x \\ \sigma x & 0 \end{bmatrix} - \begin{bmatrix} 0 & \sigma x \\ \sigma x & 0 \end{bmatrix} \begin{bmatrix} \sigma x & 0 \\ 0 & \sigma x \end{bmatrix}$$

$$= \begin{bmatrix} 0 & \sigma x^2 \\ \sigma x^2 & 0 \end{bmatrix} - \begin{bmatrix} 0 & \sigma x^2 \\ \sigma x^2 & 0 \end{bmatrix} = 0$$

From equation (7)

$$i\hbar \frac{d\sigma x'}{dt} = [\sigma x', c\alpha_y p_y + c\alpha_z p_z] = [\sigma x', c\alpha_y p_y] + [\sigma x', c\alpha_z p_z] \quad (8)$$

$$[\sigma x', \alpha_y] = \sigma x' \alpha_y - \alpha_y \sigma x'$$

$$= \begin{bmatrix} \sigma x & 0 \\ 0 & \sigma x \end{bmatrix} \begin{bmatrix} 0 & \tau y \\ \tau y & 0 \end{bmatrix} - \begin{bmatrix} 0 & \sigma y \\ \sigma y & 0 \end{bmatrix} \begin{bmatrix} \sigma x & 0 \\ 0 & \sigma x \end{bmatrix}$$

$$= \begin{bmatrix} 0 & i\sigma z \\ i\sigma z & 0 \end{bmatrix} - \begin{bmatrix} 0 & -i\tau z \\ -i\tau z & 0 \end{bmatrix} = 2i\alpha_z$$

$$\text{Similarly } [\sigma x', \sigma z] = -2i\alpha_y$$

From equation (8)

$$i\hbar \frac{d\sigma x'}{dt} = c 2i\alpha_z p_y + c(-2i\alpha_y) p_z \quad (9)$$

$$+ 2ic(\alpha_z p_y - \alpha_y p_z)$$

(5)

Multiplying both sides by $\frac{1}{2} \hbar$ to

$$i\hbar \frac{d}{dt} \left(\frac{1}{2} \hbar \sigma x' \right) = i\hbar c (\alpha_2 p_2 - \alpha_3 p_3) \quad \text{--- (10)}$$

From equation 6 & 10

$$i\hbar \frac{dLx}{dt} + i\hbar \frac{d}{dt} \left(\frac{1}{2} \hbar \sigma x' \right) = 0$$

$$= \frac{d}{dt} \left(Lx + \frac{1}{2} \hbar \sigma x' \right) = 0$$

$$Lx + \frac{1}{2} \hbar \sigma x' = \text{constant} \quad \text{--- (11)}$$

Hence the quantity $J = [L + \frac{1}{2} \hbar \vec{\sigma}]$
 commutes with H can be taken as total angular
 momentum then $S = \frac{1}{2} \hbar \vec{\sigma}$

as the spin angular momentum of a electron.
 Hence Dirac's theory automatically endows the
 spin motion of electron.